I. INTRODUCTION

Examples of “single-molecule magnets” (SMM’s) which have attracted considerable recent interest include the uniaxial \( \text{Mn}_{12}\text{O}_{12}(\text{CH}_3\text{COO})_{16}(\text{H}_2\text{O})_{4} \cdot 2\text{CH}_3\text{COOH} \cdot 4\text{H}_2\text{O} \) (Mn\(_{12}\)-acetate or Mn\(_{12}\)-ac) system and the biaxial \( \text{Fe}_8\text{O}_2(\text{OH})_{12}[(\text{tacn})_6]\text{Br}_8 \cdot 9\text{H}_2\text{O} \) (Fe\(_8\)-Br) system.\(^1\)–\(^10\) Both of these SMM’s exhibit spin \( S = 10 \) ground states, with a dominant easy-axis type anisotropy \( D S^z \) (where \( D \) is negative—see Sec. II), which provides a potential barrier (height \( \sim |D|S^2 \)) against the reversal of individual spins from \( M_S = +10 \) to \(-10 \) (see Fig. 1). Recent studies have shown that both systems possess an ability for their magnetic moments to quantum tunnel through their respective anisotropy barriers (see Fig. 1) and that this process dominates the magnetization dynamics at low temperatures (\( k_B T \ll |D|S^2 \)).\(^5\)–\(^18\)

Because of its significance as a fundamentally novel phenomenon, as well as its potential for applications in nanotechnology, magnetic quantum tunneling (MQT) in SMM’s has been extensively studied lately (see, e.g., Refs. 5 and 11–18), yet its underlying mechanism is still not fully understood, especially at a quantitative level. One possible reason is the lack of precise information on the disposition of the individual (electronic) spin energy levels. This is due to the fact that much of the available data on the spin properties of these systems have been obtained through dc and ac magnetization studies,\(^14\)–\(^16\),\(^19\) specific heat,\(^20\) neutron scattering,\(^21\)–\(^23\) proton NMR,\(^24\),\(^25\) and muon spin relaxation,\(^24\) which cannot easily or directly probe the properties of specific spin quantum levels.

An important breakthrough in the spectroscopy of SMM’s has been the development of a sensitive cavity perturbation technique for carrying out multi-high-frequency (40 to > 200 GHz) electron paramagnetic resonance (EPR) measurements on single crystals.\(^26\)–\(^30\) In particular, this technique leads to at least a three orders of magnitude improvement in detection sensitivity\(^28\) relative to conventional single-pass multi-high-frequency EPR techniques,\(^8,31–36\) which opens up many new avenues for probing SMM’s. To date, our single-crystal investigations have focused primarily on the energy dependence of the positions in the magnetic field of EPR absorptions corresponding to transitions between different spin quantum levels.\(^26\)–\(^29\) In this article, we turn our attention to the magnetic field, frequency, temperature, and spin quantum level \( (M_S) \) dependence of the EPR linewidths and, in particular, the line shapes obtained from measurements on Fe\(_8\)-Br and Mn\(_{12}\)-ac. We note that a brief account has been presented elsewhere.\(^37\) However, this earlier report emphasizes the theoretical foundations of the linewidth analysis and does not consider line shapes.

Documented lack of resolution, sensitivity, and/or various experimental artifacts have been known to affect most high-frequency EPR investigations of Fe\(_8\)-Br and Mn\(_{12}\)-ac to date. It is commonly stated that observed EPR line shapes cannot be analyzed because they contain uncontrollable distortions due to instrumental artifacts resulting from standing waves in...
the sample probe, which cause a mixing of the dissipative and reactive sample response. This is most pronounced in the case of the high-magnetic-field studies of Barra et al., asymmetric line shapes are generally observed, making it extremely difficult to evaluate the true resonance positions. Indeed, even for aligned polycrystals, significant discrepancies can be noted between the observed and calculated EPR transitions, with respect to both the peak positions and intensities, as well as the line shapes and widths. As an example, it has been reported in Ref. 33 that the linewidths in Fe₈Br approximate scale as \( M_z^2 \), whereas our earlier studies have shown a dominant linear dependence of the linewidths on \( M_z \), both for Fe₈Br and Mn₁₂-ac. This fact, along with the asymmetric line shapes observed in high-field EPR studies, may account for differences in the spin Hamiltonian parameters deduced by single-pass (i.e., cavity-less) high-field EPR, as compared to other techniques. The zero-field coherent time-domain THz technique employed by Parks et al. is limited in energy resolution (resolution is a fifth of the broadest EPR linewidth), as well as sensitivity. Consequently, it is necessary to study a large sample made up from unaligned microcrystals, and it is only possible to observe a few transitions. Furthermore, experimental complications only permit a linewidth analysis; i.e., no line shape analysis is possible. To date, the most precise SMM EPR line shape analyses have been reported by Mukhin et al., for both Fe₈Br and Mn₁₂-ac using a zero-field coherent source technique. Gaussian EPR lines are observed, which are attributed to random dipolar fields in the sample. However, these studies are also limited in sensitivity. Consequently, only a few transitions are observed, even for very large polycrystalline samples. Finally, conventional X-band EPR is limited by the fact that only a few transitions close to the top of the barrier are accessible. Furthermore, little or no information concerning the spin-spin (T₂) and spin-lattice (T₁) relaxation times has been possible in X-band.

As we will show in this article, measurements on single crystals are essential, as are several other unique aspects of our technique. For example, contrary to recent assertions by Blinc et al., we are able to obviate essentially all instrumental artifacts that could influence and/or distort EPR line shapes. In addition, our ability to measure at many frequencies, temperatures, and orientations enables us to distinguish between several contributions to the EPR linewidths and line shapes. These include inter-SMM and hyperfine dipolar fields, distributions in crystal field parameters, and spin-lattice interactions. Indeed, each of these effects have been discussed in the MQT literature, e.g., dipolar fields, distributed crystal field parameters caused by dislocations, thermally assisted tunneling, the Landau-Zener effect, and hyperfine interactions. Based on our recent theoretical work, we find that a distribution in the uniaxial crystal field parameter, so-called D strain, contributes significantly to the EPR linewidths for the Fe₈Br and Mn₁₂-ac systems. In addition, we observe an appreciable inter-SMM dipolar broadening for Fe₈Br and g strain in Mn₁₂-ac. The current work focuses on the EPR line shapes. This is particularly important in view of recent theoretical work which considers local distortions in the SMM site symmetries, and their role in the MQT phenomenon.

The article is organized as follows: in the next section (Sec. II), we discuss the \( S=10 \) effective spin Hamiltonian in the context of high-frequency EPR and the MQT phenomenon; in Sec. III, we describe our experimental procedure; in Sec. IV, we separately present results for Fe₈Br (Sec. IV A) and Mn₁₂-ac (Sec. IV B); in Sec. V, we discuss the significance of our results, and we summarize and conclude our findings in Sec. VI.

II. EFFECTIVE SPIN HAMILTONIAN

The effective spin Hamiltonian for the Mn₁₂-ac and Fe₈Br systems has the form

\[
\hat{H} = D \hat{S}_z^2 + E (\hat{S}_z^2 - \hat{S}_x^2) + \mu_B \vec{B} \cdot \vec{g} \cdot \hat{S} + \hat{O}_4 + \hat{H}',
\]

where \( \hat{S} \) is the vector spin operator and \( \hat{S}_z, \hat{S}_x, \) and \( \hat{S}_y \) are the spin projection operators along the easy, hard, and inter-
mediate axes, respectively; \(D\) is the uniaxial spin-spin coupling parameter, and the rhombic \(E(<D/3)\) term characterizes the magnetic anisotropy in the plane perpendicular to the easy axis; \(g\) is the Landé \(g\) tensor, and \(B\) is the applied magnetic field vector; finally, \(\dot{Q}_d\) denotes weaker higher-order terms of order 4 in the spin operators, and \(H'\) represents additional perturbations such as those lead to EPR line broadening, as discussed in the Introduction. \(D\) is negative for both \(\text{Mn}_{12}-\text{ac}\) [\(-0.457\) \(\text{cm}^{-1}\) (Ref. 22)] and \(\text{Fe}_8\text{Br}\) [\(-0.203\) \(\text{cm}^{-1}\) (Ref. 23)]. The rhombic term \(E\) is zero for \(\text{Mn}_{12}-\text{ac}\), due to its strictly axial (\(S_1\)) site symmetry, while \(E=0.032\) \(\text{cm}^{-1}\) for \(\text{Fe}_8\text{Br}^{23}\). For a cylindrically symmetric system, for which \(D\dot{S}^2\) would be the only nonzero term in Eq. (1) in zero applied field, the energy eigenstates may be labeled by the quantum number \(M_S(-S<M_S<S)\), which represents the projection of \(S\) onto the easy axis; the energy eigenvalues are then given by the expression \(\epsilon = DM_S^2\). This results in the energy barrier separating doubly degenerate \((M_S=\pm i,i=\text{integer})\) spin-“up” and -“down” states (see Fig. 1).

In order for quantum tunneling to be possible, there must exist finite terms in Eq. (1) that break the cylindrical symmetry (i.e., do not commute with \(\dot{S}_z\)) and mix states in either potential well, thereby lifting the \(M_S\) degeneracies. Such terms include the rhombic \(E(\dot{S}_z^2-\dot{S}_y^2)\) term, higher-order single-ion anisotropy terms in the transverse spin operators (e.g., \(\dot{S}_z^4\) and \(\dot{S}_y^4\)), \(2,4\) a transverse externally applied magnetic field, transverse internal fields due to neighboring SMM’s or nuclei (contained in \(H'\) (Refs. 39 and 46)), or local distortions in the crystal field symmetry caused imperfections in the crystal [also contained in \(H'\) (Ref. 40)]— these may include dislocations or a partial loss or disorder among the ligand molecules. Pure quantum tunneling is observed between the unperturbed \(M_S=\pm 10\) degenerate ground states in \(\text{Fe}_8\text{Br}\) below about 350 mK (see Fig. 1), and the observed relaxation has been explained in terms of the \(E(\dot{S}_z^2-\dot{S}_y^2)\) term in Eq. (1).\(^{18}\) However, one has to additionally consider the respective roles of inter-SMM dipolar fields and nuclear hyperfine fields \([H'\text{ in Eq. (1)}]\) in order to fully account for the observed tunneling in \(\text{Fe}_8\text{Br}^{23,46}\). For the \(\text{Mn}_{12}-\text{ac}\) system, the situation is even more complex: pure ground-state quantum tunneling has not been observed in zero applied magnetic field, and there is little consensus as to the origin of the apparent (weak) transverse anisotropy which causes thermally assisted quantum tunneling among levels near the top of the barrier (see Fig. 1). The leading theories include fourth-order single-ion transverse anisotropy terms such as \(\dot{S}_z^4\) and \(\dot{S}_y^4\), which are allowed under the axial (\(S_2\)) site symmetry\(^{2,4,16,22}\) and local distortions in the axial crystal fields caused, e.g., by dislocations.\(^{40}\)

A tunneling matrix element [term in Eq. (1) that does not commute with \(\dot{S}_z\)] gives rise to eigenstates that are constructed from symmetric and antisymmetric combinations of unperturbed \((M_S=\pm i)\) states on either side of the barrier and to a lifting of the degeneracy (tunnel splitting) between these symmetric and antisymmetric states. Thus, the tunneling matrix elements modify the energy eigenvalues and should, therefore, be detectable in an EPR experiment. However, for the systems of interest, these effects are small. Consequently, great precision is required. Although inelastic neutron scattering (INS) measurements are capable of providing similar information, EPR offers many advantages in terms of resolution, sensitivity (amount of sample required) and, in particular, the fact that one can systematically control the level mixing by means of an externally applied magnetic field. In addition, the higher sensitivity of the EPR technique ensures that all of this may be achieved for a single crystal. Several EPR (and INS) experiments have indicated the presence of significant fourth-order single-ion transverse anisotropies for both \(\text{Mn}_{12}-\text{ac}\) (Refs. 16, 26, 32 and 35) and \(\text{Fe}_8\text{Br}\) (Refs. 33 and 35). Nevertheless, fourth-order anisotropy \(\text{cannot}\) fully explain the observed spacing in magnetic field of magnetization steps observed in low-temperature hysteresis experiments,\(^{49}\) and there remains considerable disagreement as to the magnitude of the fourth-order terms.\(^{16,26,32–35,50}\) Consequently, higher-resolution EPR measurements are necessary in order to refine estimates of the Hamiltonian parameters and to determine whether EPR line shapes contain additional information concerning the MQT phenomenon.

### III. EXPERIMENT

The \((2S+1)\)-fold quantum energy level structure associated with a large molecular spin \(S\) necessitates spectroscopies spanning a wide frequency range. Furthermore, large zero-field level splittings, due to the significant crystalline anisotropy (large \(|D|\) and large total spin \(S\), demand the use of frequencies and magnetic fields considerably higher (50 GHz to 1 THz, and up to 10 T, respectively) than those typically used by the majority of EPR spectroscopists. The high degree of sensitivity required for single-crystal measurements is achieved using a resonant cavity perturbation technique in combination with a broadband millimeter-wave vector network analyzer (MVNA) exhibiting an exceptionally good signal-to-noise ratio; a detailed description of this instrumentation can be found in Refs. 28 and 51. The MVNA is a phase-sensitive, fully sweepable [8—350 GHz (Ref. 52)], superheterodyne source-detection system. Several sample probes couple the network analyzer to a range of high-sensitivity cavities \((Q\text{ factors of up to 25}000\text{ })\) situated within the bore of a 7 T superconducting magnet. The MVNA-cavity combination has been shown to exhibit a sensitivity of at least \(10^6\) spins \(\text{G}^{-1}\) \(\text{s}^{-1}\), which is comparable with the best narrow-band EPR spectrometers. This, coupled with newly acquired sources and a split-pair magnet, allows single-crystal measurements at any frequency in the range from 8 to 250 GHz at temperatures down to 1.5 K (\(\pm 0.01\) K) and for any geometrical combination of dc and ac field orientations up to 7 T (up to 45 T at the National High Magnetic Field Laboratory).

The use of a narrow-band cavity offers many important advantages over nonresonant methods (see, e.g., Refs. 8 and 31–34). Careful consideration concerning the coupling of radiation to and from the cavity (via waveguide), combined
with the ability to study very small samples, eliminates problems associated with standing waves in the sample probe.\textsuperscript{51} This, in turn, eliminates a mixing of the dissipative and reactive responses of the sample under investigation and, when combined with a vector detection scheme, enables faithful extraction of the true EPR line shapes (both the real and imaginary components), free from instrumental artifacts. One other notable feature of the superheterodyne scheme is its detection rate of 34 MHz; thus, there is no need for field modulation, unlike other high-field EPR methods.\textsuperscript{8,31–33} Consequently, any raw data displayed in this paper constitute pure absorption. Finally, the use of a cavity enables positioning of a single-crystal sample into a well-defined electromagnetic field environment; i.e., the orientations of the dc and ac magnetic fields relative to the sample’s crystallographic axes is precisely known, and the dc field is very homogeneous over the tiny volume ($\sim 1 \text{ mm}^3$) of the sample. In this way, additional contributions to the EPR line shapes are avoided due to, e.g., field inhomogeneities over the sample volume or slight misalignments of crystallites, as may be the case for measurements on aligned powders.

Single Mn$_{12}$-ac and Fe$_8$Br crystals were grown using literature methods.\textsuperscript{2,7} All measurements were performed in the standard EPR configuration with the ac excitation field transverse to the dc field. The Fe$_8$Br crystal used for the majority of the studies reported in the following section was a thin rhombic (acute angle $\sim 60^\circ$) platelet with edges of length $\sim 0.7$ mm and thickness $\sim 0.17$ mm. The Mn$_{12}$-ac crystal was needle shaped, with approximate dimensions $1 \times 0.15 \times 0.05$ mm$^3$. Orientation of the Mn$_{12}$-ac crystal was straightforward, due to the fact that its easy axis is defined by the needle axis of the sample. Alignment of the Fe$_8$Br crystal necessitated measurements at many angles.\textsuperscript{29} All of the data presented in this paper were obtained with the dc magnetic field aligned approximately parallel to the samples’ easy axes (hard-axis data are presented elsewhere\textsuperscript{26,29}). Finally, temperature control, in the range 2–50 K, was achieved using a variable-flow cryostat situated within the bore of the superconducting magnet.

IV. RESULTS

A. Fe$_8$Br

Figure 2 shows raw temperature-dependent data obtained for the Fe$_8$Br SMM at a frequency of 116.931 GHz; the field is applied along the easy axis to within an accuracy of about 1°. A series of more or less evenly spaced resonant absorptions are observed which become successively weaker with increasing field. A detailed analysis of the frequency dependence of the EPR line positions is published elsewhere.\textsuperscript{29} Here we focus instead on the linewidths and line shapes. In this frequency and field range, and for this orientation, the levels involved in the transitions are well separated and do not cross. Thus, $M_S$ is a good quantum number, and we have labeled the resonances in Fig. 2 according to the levels involved in the transitions.

It is apparent from Fig. 2 that the observed transitions exhibit dramatic variations in linewidth, both as a function of $M_S$, field, and temperature. In order to characterize this width dependence, it is first necessary to fit the line shapes. Figure 3 shows several attempts to fit various transitions with both Lorentzian and Gaussian functions; all of the data were obtained from a single field sweep at a frequency of 89.035 GHz and a temperature of 10 K, with the field close to alignment with the easy axis ($8^\circ \pm 1^\circ$). At the lowest fields, where the resonances are broad (comparable to spacings) and therefore overlap, fits involving multiple Gaussian or Lorentzian functions are necessary (Fig. 3). As can be seen from this figure, the Gaussian fit is noticeably better than the Lorentzian one, suggesting that some form of inhomogeneous broadening mechanism is at play. Although this has been noted by other authors,\textsuperscript{35} these earlier measurements were performed on polycrystalline powders, which were pressed into large pellets. Consequently, the microcrystals in the pellet will have been subjected to varying stresses that could be responsible for the broad Gaussian lines. Here, we demonstrate that this is an intrinsic feature of the Fe$_8$ SMM’s (this is also the case for Mn$_{12}$—see following section).

Figure 4 shows Lorentzian and Gaussian fits to transitions excited from successively higher $M_S$ levels (smaller absolute

FIG. 2. Temperature- (indicated in the figure) dependent microwave absorption amplitude for Fe$_8$Br, as a function of magnetic field, for measurements at 116.931 GHz; the field is applied parallel to the easy axis to within an accuracy of 1°. Sharp resonances are observed, which correspond to EPR transitions between the $2S + 1$ spin states of the system—these are labeled in the figure.

FIG. 3. A closeup of several Fe$_8$Br resonances obtained at 10 K and a frequency of 89.035 GHz, with the field applied approximately parallel to the easy axis; the levels involved in the transitions are indicated in the figure. The solid and dashed curves represent respective Gaussian and Lorentzian fits to the data.
value of $M_S$) from the same field sweep as the data in Fig. 3. The first thing to note is the dramatic line narrowing with decreasing absolute value of $M_S$. Also apparent (in spite of the reduced signal to noise) is a slight change in line shape, from approximately Gaussian to a more Lorentzian form, as the reduced signal to noise ratio decreases absolute value of $M_S$.

The first thing to note is the dramatic line narrowing with decreasing absolute value of $M_S$. Also apparent (in spite of the reduced signal to noise) is a slight change in line shape, from approximately Gaussian to a more Lorentzian form, as the lines become narrower. The significance of this finding will be discussed further below.

Figure 5 shows a compilation of Gaussian widths versus $M_S$ (the level from which the transition was excited), measured at many different frequencies; all of the data were obtained at a temperature of 10 K and with the field applied at an angle $\theta = 8^\circ \pm 1^\circ$ away from the easy axis. The pronounced (almost linear) increase in EPR linewidth with $M_S$ is due to the $D$ strain.** Since the dominant contribution to the spin Hamiltonian [Eq. (1)] is proportional to $M^2_S$, energy differences (i.e., EPR transition frequencies) scale as $M_S$; hence $D$ strain produces a linear width ($\Delta$ energy) dependence on $M_S$, which projects onto a width in field ($\Delta$B) that also scales linearly with $M_S$. The rounding close to $M_S=0$ is due to a convolution of the intrinsic lifetime (spin-lattice) broadening and the $M_S$-dependent ($D$-strain) contribution. The former is expected to produce a Lorentzian line shape, thus accounting for the apparent crossover from Gaussian to Lorentzian EPR lines, as observed for the sharper transitions in Fig. 4. The slight narrowing of the EPR widths in Fig. 5 as a function of increasing frequency and the weak asymmetry about $M_S = -1/2$ are due to the fact that higher-frequency (higher-$M_S$) transitions are observed at higher magnetic fields, where the inter-SMM dipolar broadening is weaker (see discussion below). For Fe$_8$Br, no measurable $g$ strain is observed. A theoretical analysis of these trends may be found in Ref. 37.

The temperature-dependent contribution to the linewidths for Fe$_8$Br (Fig. 6) is dominated by random-fluctuating dipolar fields at each SMM site, due to the magnetic moments of the neighboring SMM’s. At 116.931 GHz, the $M_S = -10$ to $-9$ transition occurs very close to zero field ($\sim 0.1$ T). Under zero-field cooling, the spin-up ($M_S = -10$) and spin-down ($M_S = +10$) states will be equally populated. Hence, the random distribution of spin-up and spin-down molecules results in a maximal contribution to the Gaussian dipolar EPR line broadening, as evidenced by the extremely large temperature-dependent contribution to the $M_S = -10$ to $-9$ transition in Fig. 6 as $T \rightarrow 0$. Upon raising the temperature, higher levels in the scheme ($|M_S| < 10$) become populated and the individual moments begin to fluctuate more rapidly, thus leading to a suppression of the dipolar broadening. The same qualitative picture may be applied to the other transitions ($M_S = -9$ to $-8$, etc.), provided that $k_BT$ exceeds the Zeeman splitting of the ground state. For $k_BT$ less than the ground-state Zeeman splitting, the system polarizes and the dipole fluctuations are suppressed, leading to an EPR line narrowing, as confirmed for all but the $M_S = -10$ to $-9$ transition in Fig. 6. The crossover occurs when the $M_S = +10$ state starts to depopulate, i.e., approximately when $k_BT = gS\mu_B B$. The inset to Fig. 6 shows the crossover temperature ($T_{CR}$), where the slope of the data in the main part.
of the figure changes sign, versus the average magnetic field strength at which each particular transition is observed (see Fig. 2). It is found that $T_{CR}$ indeed scales linearly with $B$. Furthermore, the slope of the fit to the data in the inset to Fig. 6 is $14.4 \pm 1$ K T$^{-1}$ = $gS\mu_B/k_B$, corresponding very well to $S=10$ and $g \approx 2$ to within the experimental error, thus supporting the picture of dipolar broadening. $T_{CR}$ for the 116.931 GHz $M_S=-10$ to $-9$ transition in Fig. 6 is below the minimum temperature of our current spectrometer ($T_{min} \approx 1.2$ K). Consequently, we were unable to verify whether this line eventually narrows when the spin system polarizes. However, higher-frequency and -field data, for which $T_{CR}$ should exceed the minimum temperature of our spectrometer, exhibit a similar line broadening for the $M_S=-10$ to $-9$ transition, at all temperatures investigated (see Fig. 8 below). At present, the origin of this effect is not fully understood.

In Fig. 7, we show the temperature dependent shifts of the EPR line positions obtained from Fig. 2. When compared to Fig. 6, it is noticeable that there is a similar crossover in the temperature dependence for transitions from levels with $|M_S|<10$. Therefore, this low-temperature limiting behavior may also be attributed to the polarization of the spin system. The sign of the shift indicates that the field at each SMM site, due to neighboring SMM’s, opposes the applied field; i.e., it is opposite to the polarization direction. This is to be expected for the Fe$_8$Br structure, in which the Fe$_8$ molecules reside on the corners of a slightly distorted cube. From the magnitudes of the $|M_S| < 10$ line shifts, we estimate dipolar fields at each SMM site of a few hundred gauss ($\sim 300$ G). The origin of the continued shift of the $M_S=-10$ to $-9$ transition to lower magnetic fields is, at present, not well understood.

A comprehensive theoretical analysis of this linewidth data has been published elsewhere. From such an analysis for the sample used in obtaining the data in Fig. 2, we estimate a distribution in $D$ of $\sigma_D = 0.01 D$, negligible $g$ strain, and a dipolar contribution to the broadening which assumes a nearest-neighbor SMM separation of 12 $\text{Å}$, which is in agreement with the accepted range of near-neighbor distances of between 10.5 and 14 $\text{Å}$. Within acceptable ranges, no other combination of parameter values is able to produce the same quality of fit simultaneously to all of the frequencies and temperatures studied for this particular sample, even though other parameter values may produce good fits for a particular measurement. This emphasizes the power of our multifrequency single-crystal EPR technique.

Very recent studies on a different sample from the one used in the measurements described above (similar shape and size, and obtained from the same synthesis) reveal unusual line shape characteristics, particularly at low temperatures and for transitions involving levels close to the bottom of the potential wells (e.g., $M_S=10\rightarrow -9$ and $M_S=-9 \rightarrow -8, \ldots$). This is illustrated in Fig. 8, for data obtained at a frequency of 145.937 GHz, with the field aligned precisely along the easy axis. At the somewhat elevated temperature of 10 K [Fig. 8(a)], the $M_S=-10\rightarrow -9$ transition shows appreciable asymmetry: this resonance has a Lorentzian (dotted curve) tail on the low-field side and closer to a Gaussian (dashed curve) tail on the high-field side; furthermore, the data lie somewhat to the right of the fits on the high-field side. Subsequent resonances ($M_S=-9\rightarrow -8$, $M_S=-8 \rightarrow -7$, etc.) appear to be less asymmetric, and the Gaussian fit improves. Recent theoretical studies have considered, in some detail, the effects of different types of defect structure on EPR line shapes and how these defects may influence the MQT phenomenon in SMM’s (see Sec. V). Both Lorentzian and Gaussian broadening are predicted, as well as asymmetry (see discussion in Sec. V).

Examination of the $M_S=-10\rightarrow -9$ transition at lower temperatures [Fig. 8(b)] reveals a fine structure to the line, which is strongly temperature dependent; i.e., the curves are not smooth, but exhibit a series of distinct distortions. We suspect that this fine structure and the asymmetry are connected with crystal defects, since it is not observed for all
Nevertheless, the $M_S$ would lead to a degradation of sample quality over time and metries, as opposed to a smooth distribution. The dislocation structure is indicative of a finite number of SMM site sym-

...strains in the crystal, thus giving rise to a smooth distribution of the system. The solid curves are Gaussian fits to the data; the inset shows attempts to fit both Gaussian (solid curve) and Lorentzian (dashed curve) functions to some of the data in the main part of the figure.

samples. A direct comparison between the data in Fig. 8 and data for the first sample (Figs. 3 and 4) is complicated by the fact that the linewidths are substantially broader for this second sample (increased $D$ strain). We note that many SMM compounds are known to lose solvent molecules, which would lead to a degradation of sample quality over time and may, in fact, provide a natural explanation for the $D$ strain. Nevertheless, the $M_S = -10 \to -9$ transition in Fig. 2 (lowest-field resonance) also exhibits fine structure at the lowest temperatures, and fine structure is also apparent in data published by Barra et al. It is possible that the fine structure is indicative of a finite number of SMM site symmetries, as opposed to a smooth distribution. The dislocation theory of Garanin and Chudnovsky leads to long-range strains in the crystal, thus giving rise to a smooth distribution in site symmetries. Therefore, the fine structure we observe may signify a finite number of configurations involving Fe$_8$ molecules coordinated by different numbers or arrangements of ligand molecules. It should be stressed that the data presented in Fig. 8 truly represent the intrinsic EPR line shapes; i.e., any instrumental effects would be systematic and, therefore, influence all of the line shapes equivalently. Furthermore, the noise level for these measurements is well below the resolution of the figure. Clearly, additional studies on a range of samples prepared in different ways, i.e., different sample qualities, will be required to fully resolve these issues—such studies are currently in progress.

B. Mn$_{12}$-ac

Figure 9 shows raw data obtained for Mn$_{12}$-ac at a frequency of 189.123 GHz and at temperatures between 10 and 35 K; the field is applied parallel to the easy axis to within 0.2°. The first point to note is the diminished signal-to-noise for this data, when compared to Fig. 2. There are several reasons for this: first, the dynamic range of our spectrometer is reduced at this higher frequency of 189 GHz; second, the volume of the sample used in this investigation was somewhat smaller than that of the Fe$_8$Br sample used for the studies in the previous section; finally, due to the increased energy scale of the Mn$_{12}$-ac system (larger $|D|$), the observed transitions correspond to levels considerably higher in energy above the ground state than would be the case for Fe$_8$Br at comparable frequencies and fields. Therefore, this constitutes a much more challenging experiment. Nevertheless, clear resonances are observed, which also have Gaussian line shapes (solid curves) and once again exhibit a pronounced linewidth dependence on $M_S$ (the level from which the transition was excited). For comparison, the inset to Fig. 9 shows attempts to fit the first two transitions with Gaussian and Lorentzian functions; the Gaussian fit is noticeably better.

Figure 10 plots the $M_S$ dependence of the Gaussian EPR linewidths for measurements made at several different frequencies. The linear $M_S$ contribution to the widths once again indicates a significant distribution in $D$ ($\sigma_D \approx 0.02D$). However, the dipolar contribution to the linewidths is considerably weaker for Mn$_{12}$-ac than for Fe$_8$Br. This is initially apparent from the weak frequency dependence observed among the widths for transitions with $|M_S| < 5$ and the fact that the frequency dependence observed for higher-$|M_S|$ transitions is opposite to what was observed for Fe$_8$ in Fig. 5; i.e., the higher-frequency transitions are broader. The origin of this effect has to do with $g$ strain and will be discussed further below.

In Fig. 11, we plot the temperature dependence of the linewidths observed at a frequency of 189.123 GHz. A similar trend can be observed for the low-$|M_S|$ transitions for both Mn$_{12}$-ac and Fe$_8$Br (decreasing width with decreasing temperature). However, the pronounced line narrowing observed for the higher-$|M_S|$ transitions at lower temperatures, in Fe$_8$Br, is absent in Mn$_{12}$-ac. This provides further evidence that the dipolar contribution to the linewidths is considerably weaker in Mn$_{12}$-ac. In contrast, the $M_S$ dependence of the widths is stronger for Mn$_{12}$-ac, due to the increased $D$ strain ($\sigma_D \approx 0.02D$, as opposed to $\sigma_D \approx 0.01D$ for Fe$_8$Br). The weaker dipolar broadening is consistent with an increased inter-SMM separation for Mn$_{12}$-ac.
The pronounced frequency dependence of the larger $|M_S|$ ($>5$) transitions, as seen in Fig. 10, is attributable to g strain. Each transition shifts to higher field with increasing frequency, thereby amplifying the effect of a distribution in $g$—hence the increased linewidths with increasing frequency. Additional support for $g$ strain comes from the measured anisotropy in the $g$ tensor, from both our previous work \cite{note1} and that of others. \cite{note2,note3} Using accepted values for the inter-SMM separations, it is possible to simulate all data obtained for a particular sample with a single pair of (sample-dependent) $D$- and $g$-strain parameters; $\sigma_D = 0.02D$ and $\sigma_g = 0.008g$ are typical for many of the Mn$_{12}$-ac samples we have studied. \cite{note4}

Finally, in Fig. 12, we show magnifications of several of the resonances from Fig. 9. Once again, it is apparent that the EPR lines exhibit a fine structure. In particular, the $M_S = -6$ to $-5$ and $M_S = -5$ to $-4$ transitions show a clear double-peak structure; the best attempts at single-Gaussian fits (dotted curves) are included in the figure for comparison. Similar behavior was observed at other frequencies. Once again, these findings hint at the possibility that there may exist a finite number of configurations involving Mn$_{12}$-ac molecules coordinated by different numbers or arrangements of ligand molecules. We note that the presence of various isomeric forms of Mn$_{12}$, whereby the arrangements of the ligands differ slightly, is well established both in the acetate \cite{note5} and in Mn$_{12}$ SMM’s involving other ligands. \cite{note6}


V. DISCUSSION

The unambiguous demonstration of $D$ strain in Fe$_8$Br and Mn$_{12}$-ac indicates the presence of significant sample inhomogeneities or imperfections; i.e., the molecules are subjected to varying anisotropy. A natural explanation for this could be a loss of solvent molecules or disorder among the ligands; indeed, a slight degradation of sample quality over time has been noted, and it is known that the crystals do lose solvent. \cite{note7} Another possibility has recently been proposed by Garanin and Chudnovsky, \cite{note8} involving dislocations in the crystals which act to locally distort the site symmetries of individual SMM’s. What is more, Garanin and Chudnovsky have claimed that the QTM phenomenon in SMM’s finds a natural explanation in terms of these dislocations. As mentioned in the Introduction, QTM requires symmetry breaking terms in the Hamiltonian [Eq. (1)]. Local rotations of the magnetic anisotropy axes due to dislocations have two effects: they generate transverse second-order crystal field terms, of the form $E_i (\vec{S}_z^2 - \vec{S}_y^2)$, which facilitate QTM between $M_S$ states differing by an even integer, and they result in an effective local transverse magnetic field (under application of field $||z$) that unfreezes odd tunneling resonances, thereby solving the conundrum as to why both even and odd tunneling resonances are always observed in Fe$_8$Br and Mn$_{12}$-ac. It has also been shown that long-range strains produced by dislocations result in broad distributions of relaxation times, something which is also observed experimentally in Mn$_{12}$-ac. \cite{note9} Indeed, assuming realistic concentrations of dislocations, Garanin and Chudnovsky \cite{note9} have been able to account for the experimentally observed tunneling in Mn$_{12}$-ac (Ref. 49) and predict a distribution in the uniaxial spin-spin coupling parameter $\sigma_p \approx 0.027D$, which is of the same order as we have found from linewidth analyses [$\sigma_D \approx (0.01–0.02)D$]. It should be pointed out, however, that any disorder amongst the ligand molecules, as well as any loss of solvent molecules, may be expected to produce many of the same effects as those proposed by Garanin and Chudnovsky; \cite{note9} i.e., disorder among the peripheral molecules surrounding the SMM’s will result in local variations in site symmetry.

Careful theoretical consideration of different forms of dislocations (edge, screw, correlated, uncorrelated, etc.) indicates that the resultant $D$ strain can produce both Gaussian and Lorentzian EPR line shapes. \cite{note10} Furthermore, recent work by Park et al. \cite{note10} predicts asymmetry in some of the EPR line shapes. One may think of this as arising due to the fact that a dislocation causes an asymmetric distribution in $D$; i.e., it can only reduce $D$, not enhance it. Clearly, certain features of the results presented here, particularly in Fig. 8(a), may confirm some of these predictions. So far, our attempts to simulate data have been limited to $D$ strain, $g$ strain, and dipolar broadening. The $D$ and $g$ strains should be correlated, although we have yet to verify this. Inclusion of $E$ strain, as well as strain in the fourth-order anisotropy, did not affect
VI. SUMMARY AND CONCLUSIONS

We have used a multi-high-frequency cavity perturbation technique to measure EPR line shapes for oriented single-crystal samples of the Fe$_8$Br and Mn$_{12}$-ac SMM's. The use of a narrow-band cavity offers many important advantages over nonresonant methods: careful consideration concerning the coupling of radiation to and from the cavity (via waveguide), combined with the ability to study very small samples, eliminates mixing of the dissipative and reactive responses of the sample under investigation, and enables faithful extraction of the true EPR line shapes. Our ability to measure at many frequencies, temperatures, and orientations further allows us to distinguish between several contributions to the EPR linewidths and shapes. These include inter-SMM and hyperfine dipolar fields, distributions in crystal field parameters, and spin-lattice interactions. Each of these effects have been discussed in the literature in terms of playing a possible role in the MQT phenomenon. In addition to significant $D$ strain, which is indicative of a distribution in the SMM site symmetries, we have detected a previously unobserved fine structure to some of the resonances for both Fe$_8$Br and Mn$_{12}$-ac SMM's.

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The identical experimental setup as that used in Ref. 26 has been used for single-crystal EPR by E. del Barco, J.M. Hernandez, J. A.A.J. Perenboom, M. R. Schaapman, S. O. Hill, and J. S. Brooks, Polyhedron, 16, 237 (1999).

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This frequency coverage will soon be extended to $>600$ GHz.


The dipolar calculation makes the rather bold simplification of assuming the SMM’s to be distributed on a simple cubic lattice.

Very recent measurements on this second sample were made approximately one year after the original experiments were performed on the first sample (Figs. 2–7). During this time, the samples were stored in a refrigerator ($\sim$4 °C). It is possible that the samples lost solvent molecules during this period, as evidenced by the changes in line shape which are discussed in main text of this article.

K. Park, M. A. Novotny, N. S. Dalal, S. Hill, and P. A. Rikvold, J. Appl. Phys. (to be published); also cond-mat/0111589 (unpublished); also cond-mat/0112150 (unpublished).

G. Christou (private communication).


S. Hill et al. (unpublished).